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ITERATIVE ADAPTIVE IMPLICIT–EXPLICIT METHODS FOR FLOW PROBLEMS

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SUMMARY

Iterative versions of the adaptive implicit–explicit method are presented for the finite element computation of flow problems with particular reference to incompressible flows and advection–diffusion problems. The iterative techniques employed are the grouped element-by-element and generalized minimum residual methods.

KEY WORDS Adaptive implicit–explicit Grouped element-by-element

1. INTRODUCTION

For most flow problems of practical interest, especially in three dimensions, methods involving direct solution of linear equation systems give rise to massive global coefficient matrices. Storage and inversion of these matrices place a heavy demand on the computational resources in terms of the CPU time and memory. Iterative, explicit and semi-explicit methods, on the other hand, require storage and inversion of much simpler matrices, if at all. Properly implemented, these methods can substantially reduce the demand for memory and CPU time while retaining most of the desirable properties of the direct methods. In this paper we present the adaptive implicit–explicit (AIE) procedures[1] which are employed in combination with iterative techniques such as the grouped element-by-element (GEBE)[2] and generalized minimum residual (GMRES)[3] methods. Although the description of the concepts and the numerical demonstrations are based on incompressible flows and convection–diffusion problems, the approach presented here is applicable to a wider class of problems in computational fluid dynamics.

In the AIE method the elements are dynamically grouped into implicit and explicit subsets. The selection of the implicit elements is determined, at a given instant in time, by the element level Courant number and some measure of the local variations in the solution. Since the computational cost associated with the explicit elements is much smaller than that of the implicit elements, by placing the implicit elements only where and when they are needed, substantial savings in the CPU time and memory can be achieved.

The GEBE iteration method, on the other hand, is based on static (i.e. one-time) arrangement of the elements into groups, with the condition that no two elements in the same group can share a common node. In the GEBE method the preconditioning matrix is chosen to be a sequential product of the element group matrices; this approach is a variation of the one taken in the regular EBE methods employed in computational fluid dynamics[4–6] and solid mechanics.[7] The GEBE approach eliminates the need for the formation, storage and factorization of large global matrices.

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The element level matrices can be either stored or recomputed; in the case when they are stored, the storage needed is still only linearly proportional to the number of elements. To minimize the overhead associated with the synchronization involved in moving from one group to another, we attempt to minimize the number of groups. Furthermore, to increase the vector efficiency of the computations performed, within each group the elements are processed in packets of 128 elements.

The GMRES method was proposed by Saad and Schultz.[3] It is based on the minimization of the residual norm over a Krylov space. This method, with a properly chosen preconditioner, can improve the convergence rate of the iterative algorithms for non-symmetric systems substantially. Applications of the GMRES method to various fluid dynamics problems, including compressible and incompressible flows, can be found in References 6 and 8.

In the iterative versions of the AIE method, to solve the equation systems resulting from the implicitly treated elements, we employ the GEBE or GMRES iteration method. This approach leads to an iterative AIE scheme which involves no direct solution effort. In the limiting case, if we choose to treat all elements implicitly then the method becomes a pure GEBE or GMRES iteration method.

2. THE FORMULATION

Consider a two-dimensional spatial domain Ω and a time interval $(0, T)$ with $\mathbf{x}$ and t representing the co-ordinates associated with Ω and $(0, T)$. In two-dimensional space the vorticity–streamfunction formulation of the incompressible Navier–Stokes equations consists of a time-dependent advection–diffusion equation for the vorticity ω,

$$\partial\omega/\partial t + \mathbf{u}\cdot\nabla\omega - \nu\nabla^2\omega = 0 \quad \text{on } \Omega\times(0, T), \tag{1}$$

and a Poisson equation for the streamfunction ψ,

$$\nabla^2\psi + \omega = 0 \quad \text{on } \Omega\times(0, T), \tag{2}$$

where

$$\mathbf{u} = \{\partial\psi/\partial x_2, -\partial\psi/\partial x_1\} \tag{3}$$

is the velocity and ν is the kinematic viscosity. The boundary conditions associated with (1) and (2) are rather involved; we refer the interested reader to Reference 9. We note that the convection–diffusion equation governing the transport of a passive contaminant is a special case of (1) in which the velocity field is known.

A proper finite element formulation of the problem results in the following equation system for the incremental values of the nodal unknown vectors:

$$\begin{vmatrix} \mathbf{A}_{11} & \mathbf{A}_{12} & \mathbf{A}_{13} \\ \mathbf{A}_{21} & \mathbf{A}_{22} & \mathbf{A}_{23} \\ \mathbf{A}_{31} & \mathbf{A}_{32} & \mathbf{A}_{33} \end{vmatrix} \begin{vmatrix} \Delta\boldsymbol{\omega}_* \\ \Delta\boldsymbol{\psi} \\ \Delta\boldsymbol{\omega}_G \end{vmatrix} = \begin{vmatrix} \mathbf{R}_1 \\ \mathbf{R}_2 \\ \mathbf{R}_3 \end{vmatrix}. \tag{4}$$

The vectors $\boldsymbol{\psi}, \boldsymbol{\omega}_*$ and $\boldsymbol{\omega}_G$ represent the unknown nodal values of the streamfunction, the vorticity at the interiors and the vorticity at the boundaries respectively.

Solution of the coupled system (4) by a direct method such as Gaussian elimination places a prohibitive burden in terms of the CPU time and memory for large-scale problems. Alternatively we can consider a block-iteration scheme in which the following uncoupled equation systems are

solved iteratively until a predetermined convergence condition is met:

$$\mathbf{A}_{11}\Delta\omega_* = \mathbf{R}_1 \quad \text{(block 1)}, \tag{5}$$

$$\mathbf{A}_{22}\Delta\psi = \mathbf{R}_2 \quad \text{(block 2)}, \tag{6}$$

$$\mathbf{A}_{33}\Delta\omega_G = \mathbf{R}_3 \quad \text{(block 3)}, \tag{7}$$

Remark 1. Although under certain conditions the matrix $\mathbf{A}_{11}$ can be symmetric and positive-definite,[9] we assume that in general this is not the case.

Remark 2. $\mathbf{A}_{22}$ is symmetric and positive-definite.

Remark 3. With proper implementation, $\mathbf{A}_{33}$ can be of tridiagonal form; this makes the solution of this block essentially as easy as the solution of a one-dimensional problem.

3. THE SOLUTION TECHNIQUES

In our block-iteration procedure, at every iteration we need to solve three equation systems: (5), (6) and (7). The cost involved in (7) is quite minor (see Remark 3) and therefore we solve this equation with a direct method. Our main objective here is to minimize the computational cost associated with solving equations (5) and (6), which, after hiding all the subscripts, can be rewritten in the following general form:

$$\mathbf{A}\mathbf{x} = \mathbf{b}. \tag{8}$$

For completeness we first briefly describe the AIE, GEBE and GMRES methods.

The adaptive implicit–explicit (AIE) method

Let $\mathscr{E}$ be the set of all elements $e = 1, 2, \ldots, n_{el}$ where n_{el} is the number of elements. The assembly of the global matrix $\mathbf{A}$ can be expressed as

$$\mathbf{A} = \sum_{e\in\mathscr{E}} \mathbf{A}^e, \tag{9}$$

where $\mathbf{A}^e$ is the component of $\mathbf{A}$ contributed by element e.

The AIE method is based on partitioning of the set of elements into the subsets $\mathscr{E}_I$ and $\mathscr{E}_E$ such that

$$\mathscr{E} = \mathscr{E}_I \cup \mathscr{E}_E, \tag{10}$$

$$\varnothing = \mathscr{E}_I \cap \mathscr{E}_E. \tag{11}$$

We then replace with

$$\mathbf{A}_{AIE} = \sum_{e\in\mathscr{E}_I} \mathbf{A}^e + \sum_{e\in\mathscr{E}_E} (\mathbf{A}^e)_E, \tag{12}$$

in which

$$(\mathbf{A}^e)_E = \text{lump}(\mathbf{M}^e) \quad \text{for block 1} \tag{13}$$

and

$$(\mathbf{A}^e)_E = \text{diag}(\mathbf{A}^e) \quad \text{for block 2.} \tag{14}$$

Here lump $(\mathbf{M}^e)$ is the lumped version of the mass matrix for element e and diag $(\mathbf{A}^e)$ is the diagonal of $\mathbf{A}^e$. The matrix $\mathbf{A}_{AIE}$ has a skyline profile which is typically as shown in Figure 1.

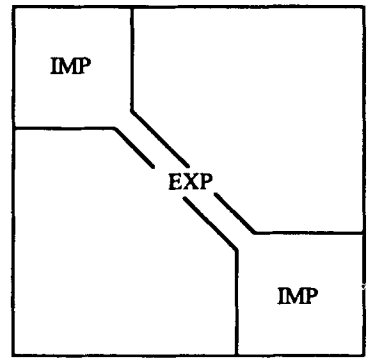

Figure 1. Typical skyline profile of the matrix $\mathbf{A}_{\mathrm{AIE}}$

In block 1 we use a direct method to solve (5), whereas in block 2 we use $\mathbf{A}_{\mathrm{AIE}}$ as the preconditioner for the conjugate gradient method[1] employed to solve (6). Because the bandwidth for the parts of $\mathbf{A}_{\mathrm{AIE}}$ corresponding to the explicit regions is substantially reduced, the method leads to savings in CPU time and memory.

The grouping given by (10) and (11) is achieved dynamically (adaptively) on the basis of element level stability and accuracy considerations.

The stability criterion is given in terms of the element Courant number $C_{\Delta t}$, which is defined as

$$C_{\Delta t} = \frac{\|\mathbf{u}\| \Delta t}{h}, \tag{15}$$

where h is the 'element length'.[10] Any element with Courant number greater than the stability limit of the explicit method needs to belong to the implicit group $\mathscr{E}_{\mathrm{I}}$.

For accuracy considerations we want to use a test parameter $\sigma^e_{\mathscr{E}}$ which is a measure of the local variations in the solution. One possible way is to define this test parameter on the basis of the element level L^2-norm of the residual r; we borrow this idea from the adaptive mesh refinement techniques given in Reference 11; that is,

$$\sigma^e_{\mathscr{E}} = \frac{\|r\|^0_{\Omega^e} - \min\limits_e (\|r\|^0_{\Omega^e})}{\max\limits_e (\|r\|^0_{\Omega^e}) - \min\limits_e (\|r\|^0_{\Omega^e})}, \tag{16}$$

where

$$\|r\|^0_{\Omega^e} = \left(\int_{\Omega^e} r^2 \, \mathrm{d}\Omega \right)^{1/2}. \tag{17}$$

Elements with $\sigma^e_{\mathscr{E}}$ greater than a predetermined value belong to group $\mathscr{E}_{\mathrm{I}}$. For other choices for $\sigma^e_{\mathscr{E}}$ see Reference 1.

Implementation of the AIE scheme is quite straightforward; compared to adaptive schemes based on grid moving or element subdividing, it involves minimal bookkeeping and no geometric constraints.

The grouped element-by-element (GEBE) method

In this method the elements are arranged into N_{pg} groups with the provision that no two elements within a group can share a common node. This way, within each group, computations

performed in element-by-element fashion can be done in parallel. In parallel computations we would like to minimize the synchronization overhead associated with finishing with one group and starting with another one. For this purpose the element-grouping algorithm described in Reference 2 tries to minimize the number of groups. We note that this grouping is a static (one-time) kind and therefore the computational cost involved in achieving it is a one-time cost. Furthermore, within each goup the elements are processed in packets of 128 (or an appropriate size) elements. This increases the vector efficiency of the computations.

On the basis of the grouping, the matrix $\mathbf{A}$ can be written as

$$\mathbf{A} = \sum_{K=1}^{N_{\mathrm{pg}}} \mathbf{A}_K, \tag{18}$$

with the 'group matrices' defined as

$$\mathbf{A}_K = \sum_{e \in \mathscr{E}_k} \mathbf{A}^e, \quad K = 1, 2, \ldots, N_{\mathrm{pg}}, \tag{19}$$

where $\mathscr{E}_K$ is the set of elements which belong to group K.

We start with the following scaled version of (8):

$$\tilde{\mathbf{A}}\tilde{\mathbf{x}} = \tilde{\mathbf{b}}, \tag{20}$$

where

$$\tilde{\mathbf{A}} = \mathbf{W}^{-1/2}\mathbf{A}\mathbf{W}^{-1/2}, \tag{21}$$

$$\tilde{\mathbf{x}} = \mathbf{W}^{1/2}\mathbf{x}, \tag{22}$$

$$\tilde{\mathbf{b}} = \mathbf{W}^{-1/2}\mathbf{b}, \tag{23}$$

and $\mathbf{W}$ is the scaling matrix. Selection of this scaling matrix depends on the properties of $\mathbf{A}$; the two choices we have considered are lump ($\mathbf{M}$) and diag ($\mathbf{A}$).

In the preconditioned iteration method, at iteration m, the following equation system is solved for $\Delta\tilde{\mathbf{y}}_m$:

$$\tilde{\mathbf{P}}\Delta\tilde{\mathbf{y}}_m = \tilde{\mathbf{r}}_m, \tag{24}$$

where $\tilde{\mathbf{P}}$ is the preconditioning matrix and the residual vector $\tilde{\mathbf{r}}_m$ is defined as

$$\tilde{\mathbf{r}}_m = \tilde{\mathbf{b}} - \tilde{\mathbf{A}}\tilde{\mathbf{x}}_m. \tag{25}$$

If $\mathbf{A}$ is symmetric and positive-definite then the vector $\tilde{\mathbf{x}}_m$ is updated by using a conjugate gradient method; otherwise we update this vector according to the expression

$$\tilde{\mathbf{x}}_{m+1} = \tilde{\mathbf{x}}_m + s\Delta\tilde{\mathbf{y}}_m, \tag{26}$$

for which the search parameter s is determined with the formula

$$s = \frac{(\tilde{\mathbf{A}}\Delta\tilde{\mathbf{y}}_m)\cdot\tilde{\mathbf{r}}_m}{\|\tilde{A}\Delta\tilde{y}_m\|^2}; \tag{27}$$

this formula is obtained by minimizing $\|\tilde{\mathbf{r}}_{m+1}\|^2$ with respect to s.

Remark 4. In the evaluation of the residual vector (25), the matrix–vector multiplication is performed in element-by-element fashion as shown below:

$$\tilde{\mathbf{A}}\tilde{\mathbf{x}} = \sum_{e=1}^{n_{el}} \tilde{\mathbf{A}}^e\tilde{\mathbf{x}}. \tag{28}$$

Therefore the residual vector computations are highly vectorizable. In our computations we choose to store the element level matrices.

In our GEBE approach, for block 1 (i.e. equation (5)) we use the two-pass GEBE preconditioner, whereas for block 2 (i.e. equation (6)) we use the GEBE preconditioner based on Crout factorization. We give a brief description of these two preconditioners.

The two-pass GEBE preconditioner (2P-GEBE). This preconditioning matrix, in its scaled form, is defined as

$$\tilde{\mathbf{P}} = \prod_{K=1}^{N_{pg}} \tilde{\mathbf{E}}_K \prod_{K=N_{pg}}^{1} \tilde{\mathbf{E}}_K, \tag{29}$$

where

$$\tilde{\mathbf{E}}_K = \mathbf{I} + \tfrac{1}{2}\tilde{\mathbf{B}}_K, \qquad K = 1, 2, \ldots, N_{pg}, \tag{30}$$

with

$$\tilde{\mathbf{B}}_K = \tilde{\mathbf{A}}_K - \tilde{\mathbf{W}}_K, \qquad K = 1, 2, \ldots, N_{pg}, \tag{31}$$

$$\tilde{\mathbf{W}}_K = \mathbf{W}^{-1/2}(\mathbf{W}_K)\mathbf{W}^{-1/2}, \quad K = 1, 2, \ldots, N_{pg}. \tag{32}$$

The definition given by (31) leads to 'Winget regularization'; we have also been experimenting with an alternative definition given as

$$\tilde{\mathbf{B}}_K = \tilde{\mathbf{A}}_K, \quad K = 1, 2, \ldots, N_{pg}. \tag{33}$$

Remark 5. Since there is no inter-element coupling within each group, $\tilde{\mathbf{E}}_K$ can also be written as

$$\tilde{\mathbf{E}}_K = \prod_{e \in \mathcal{E}_K} (\mathbf{I} + \tfrac{1}{2}\tilde{\mathbf{B}}^e), \quad K = 1, 2, \ldots, N_{pg}. \tag{34}$$

The GEBE preconditioner based on Crout factorization (Crout-GEBE). Consider the following Crout factorization:

$$\mathbf{I} + \tilde{\mathbf{B}}_K = \hat{\mathbf{L}}_K \hat{\mathbf{D}}_K \hat{\mathbf{U}}_K, \quad K = 1, 2, \ldots, N_{pg}. \tag{35}$$

The Crout-GEBE preconditioner, in its scaled form, is defined as

$$\tilde{\mathbf{P}} = \prod_{K=1}^{N_{pg}} \hat{\mathbf{L}}_K \prod_{K=1}^{N_{pg}} \hat{\mathbf{D}}_K \prod_{K=N_{pg}}^{1} \hat{\mathbf{U}}_K. \tag{36}$$

Details on vectorization and parallel processing of the GEBE method can be found in Reference 12.

The generalized minimum residual (GMRES) method

For block 1 we also have the option to employ the GMRES method; an outline of the version of the GMRES method used is given below.

Given $\mathbf{x}_0$,

set $m = 0$.

(i) Calculate the residual:

$$\mathbf{r}_m = \mathbf{W}^{-1}(\mathbf{A}\mathbf{x}_m - \mathbf{b}). \tag{37}$$

(ii) Construct the Krylov space:

$$\mathbf{e}^{(1)} = \mathbf{r}_m / \|\mathbf{r}_m\|, \tag{38}$$

$$\mathbf{f}^{(j)} = \mathbf{W}^{-1}\mathbf{A}\mathbf{e}^{(j-1)} - \sum_{i=1}^{j-1} (\mathbf{W}^{-1}\mathbf{A}\mathbf{e}^{(j-1)}, \mathbf{e}^{(i)})\mathbf{e}^{(i)}, \quad 2 \leqslant j \leqslant k, \tag{39}$$

$$\mathbf{e}^{(j)} = \mathbf{f}^{(j)} / \|\mathbf{f}^{(j)}\|, \tag{40}$$

where k is the dimension of the Krylov space.

(iii) Update the unknown vector:

$$\mathbf{x}_{m+1} = \mathbf{x}_m + \sum_{j=1}^{k} s_j \mathbf{e}^{(j)}, \tag{41}$$

where $\mathbf{s} = \{s_j\}$ is the solution of the equation system

$$\mathbf{Q}\mathbf{s} = \mathbf{z}, \tag{42}$$

with

$$\mathbf{Q} = [(\mathbf{W}^{-1}\mathbf{A}\mathbf{e}^{(i)}, \mathbf{W}^{-1}\mathbf{A}\mathbf{e}^{(j)})], \quad 1 \leqslant i, j \leqslant k, \tag{43}$$

$$\mathbf{z} = \{(\mathbf{W}^{-1}\mathbf{A}\mathbf{e}^{(i)}, -\mathbf{r}_m)\}, \quad 1 \leqslant i \leqslant k. \tag{44}$$

(iv) Go to the next iteration:

$n \leftarrow n+1$ and go to (37).

The iterations continue until $\|\mathbf{r}_m\|$ becomes less than a predetermined value. We note that the matrix $\mathbf{Q}$ is symmetric and positive-definite. Again, for the scaling matrix $\mathbf{W}$ we consider the choices lump ($\mathbf{M}$) and diag ($\mathbf{A}$). For other kinds of preconditioners see References 7 and 8.

Combinations of the AIE and iterative methods

Within the framework of our AIE scheme we can employ an iterative technique, such as the GEBE or GMRES method, to solve the equation system resulting from the implicitly treated elements. This way those elements which need to be treated implicitly are treated so, any yet the scheme involves no direct solution effort. In the rest of this section we describe two of the several possible combinations.

AIE/GEBE on block 1 and GEBE on block 2. In this method, for block 1 we use the AIE technique with the 2P-GEBE method employed to solve the equation system resulting from the implicitly treated elements; for block 2 we use the Crout-GEBE preconditioned conjugate gradient method.

The element grouping concept still applies in this method. In the implicit zones we still have groups within which the elements have no common nodal points. We do not need to redo the element grouping every time the distribution of the implicit zones is changed. The implicit elements are selected from the entire set of elements which are already grouped. The parallel nature of the GEBE method therefore is not affected by mixing with the AIE scheme.

AIE/GMRES on block 1 and GEBE on block 2. This time, for block 1 we use the AIE technique with the GMRES method employed to solve the equation system resulting from the implicitly treated elements; for block 2 we use, again, the Crout-GEBE preconditioned conjugate gradient method.

In the AIE/GMRES algorithm we initialize the unknown vector as

$$\mathbf{x}_0 = (\text{lump}(\mathbf{M}))^{-1}\mathbf{b}. \tag{45}$$

This way, for explicitly treated equations the corresponding part of $\mathbf{x}_0$ is accepted as the solution, whereas for the remaining equations the corresponding part of $\mathbf{x}_0$ is used as the initial guess for the GMRES iterations. The element-grouping concept remains in effect and facilitates the vectorization and potential parallel processing.

4. NUMERICAL EXAMPLES AND BENCHMARKING

All computations were performed on the Minnesota Supercomputer Center CRAY-2 (four CPUs, 512 megawords of memory, 4·1 ns clock and UNICOS 4·0 operating system).

Two-dimensional rotational advection of a cosine hill (the rotating puff problem)

This standard pure advection problem is used in this paper mainly for the purpose of benchmarking based on the CPU time and memory requirements. We use a unit square domain and employ uniform meshes with $15n \times 15n$ elements, where $n = 1, 2, 4, 8$ and 16. The velocity field is rotational with respect to the centre of the domain. Initially, the cosine hill has unit amplitude and a base radius of 0·2, and is centred at (0·27, 0·5). The time step for each mesh is chosen to give a Courant number of 0·22 at the peak of the cosine hill. All boundary conditions are Dirichlet-type.

The critical values for the test parameters $C_{\Delta t}$ and σ_g^e are 1·0 and 0·05 respectively. The convergence limit for the iterative solvers is set to 10^{-7}. For the GMRES method the dimension of the Krylov space is five. The results for the benchmarking based on the CPU time and memory requirements for the implicit, AIE, GEBE, GMRES, AIE/GEBE and AIE/GMRES methods are shown in Tables I and II.

It should be noted that some vectorization breakdown is involved in the selective treatment of the implicit and explicit elements; therefore an increase in CPU time, compared to the pure GEBE and GMRES methods, can be observed in some cases.

Table I. The results for the benchmarking based on the CPU time for various methods applied to the rotating puff problem

Mesh	IMP	AIE	GEBE	GMRES	AIE/GEBE	AIE/GMRES
15 × 15	1·0	0·880	1·649	1·124	1·368	1·115
30 × 30	1·0	0·722	1·153	0·937	1·039	0·924
60 × 60	1·0	0·553	0·775	0·704	0·743	0·692
120 × 120	1·0	0·385	0·430	0·425	0·437	0·423
240 × 240	1·0	0·174	0·209	0·209	0·212	0·206

Table II. The results for the benchmarking based on the memory needed for the coefficient matrices for various methods applied to the rotating puff problem

Mesh	IMP	AIE	GEBE	GMRES	AIE/GEBE	AIE/GMRES
15 × 15	1·0	0·343	1·552	1·123	0·570	0·810
30 × 30	1·0	0·214	0·726	0·537	0·243	0·366
60 × 60	1·0	0·167	0·352	0·263	0·114	0·176
120 × 120	1·0	0·142	0·173	0·130	0·055	0·087
240 × 240	1·0	0·113	0·086	0·065	0·026	0·043

Flow past a circular cylinder

For this benchmark problem we used three different finite element meshes. Mesh A consists of 1310 elements and 1365 nodes; around the cylinder there are 29 elements in the radial direction and 40 elements in the circumferential direction. Mesh B involves 5220 elements and 5329 nodes with 58 and 80 elements in the radial and circumferential directions respectively. Mesh C contains 19 836 elements and 20 046 nodes with 116 and 156 elements in the radial and circumferential directions respectively. The dimensions of the computational domain, normalized by the cylinder diameter, are 30·5 and 16·0 in the flow and cross-flow directions respectively. The free stream velocity is 0·125 and the initial value of the vorticity is zero everywhere in the domain. The Reynolds number based on the uniform free stream velocity and the cylinder diameter is 100.

The critical values for the test parameters $C_{\Delta t}$ and σ_g^e are 1·0 and 10^{-5} respectively.

For the GMRES method the dimension of the Krylov space is five. The convergence limit for the iterative solvers is 10^{-7} for block 1 and 10^{-6} for block 2. We tested seven methods: implicit, block iteration. AIE, GEBE (2P-GEBE on block 1 and Crout-GEBE on block 2), GMRES (GMRES on block 1 and Crout-GEBE on block 2), AIE/GEBE (AIE/GEBE on block 1 and Crout-GEBE on block 2) and AIE/GMRES (AIE/GMRES on block 1 and Crout-GEBE on block 2). The results for the benchmarking based on the CPU time and memory requirements are shown in Tables III and IV.

For this problem the solutions obtained with the AIE method can be found in Reference 1.

Plane jet impinging on a wedge

In this problem we illustrate how the iterative AIE method works; we employ the AIE/GEBE method (AIE/GEBE on block 1 and Crout-GEBE on block 2). The computational domain is an 80×80 square, and the distance between the jet and the leading tip of the wedge is 7·5. The single mesh employed contains 10 566 nodal points and 10 296 elements (see Figure 2). The jet inlet consists of a parabolic velocity profile with both the width and the mean value set to unity; the Reynolds number based on these values is 250. The computation is performed with a time step size 0·05. The critical values for the test parameters $C_{\Delta t}$ and σ_g^e are 1·0 and 10^{-5} respectively. The

Table III. The results for the benchmarking based on the CPU time for various methods applied to flow past a circular cylinder

Mesh	IMP	BLOCK	AIE	GEBE	GMRES	AIE/GEBE	AIE/GMRES
A	1·0	0·478	0·542	0·165	0·191	0·167	0·193
B	1·0	0·578	0·423	0·118	0·159	0·139	0·160
C	1·0	0·797	0·349	0·159	0·169	0·165	0·166

Table IV. The results for the benchmarking based on the memory needed for the coefficient matrices for various methods applied to flow past a circular cylinder

Mesh	IMP	BLOCK	AIE	GEBE	GMRES	AIE/GEBE	AIE/GMRES
A	1·0	0·399	0·084	0·172	0·148	0·104	0·123
B	1·0	0·407	0·080	0·085	0·073	0·051	0·061
C	1·0	0·410	0·077	0·042	0·037	0·025	0·030

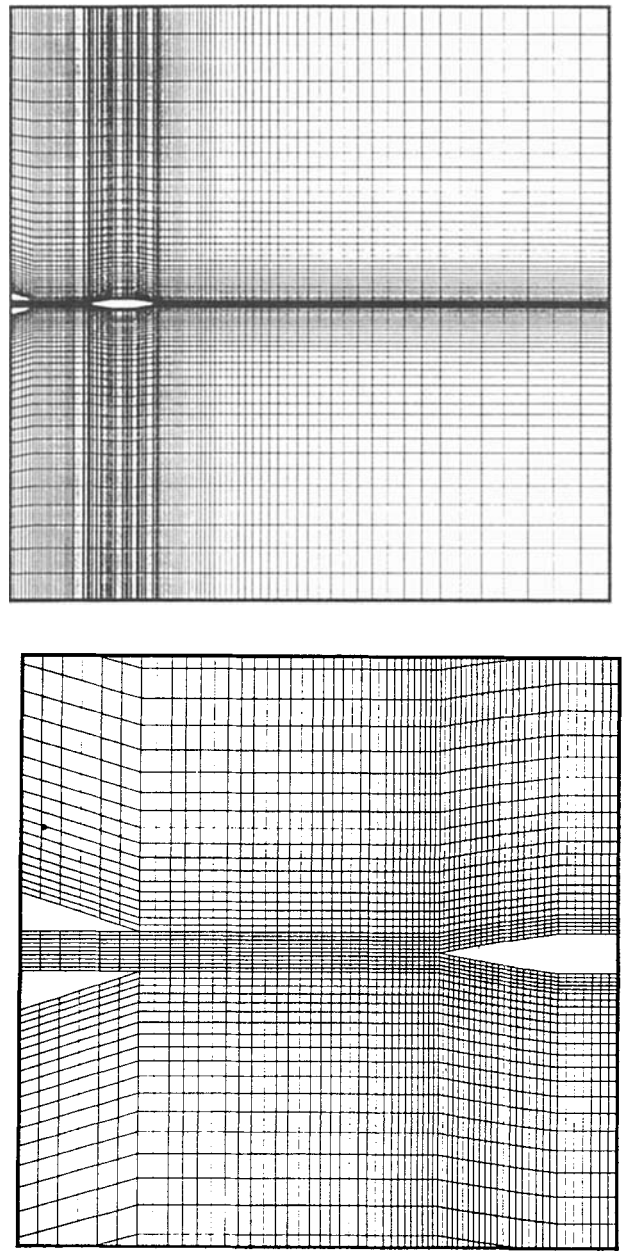

Figure 2. Plane jet impinging on a wedge: the finite element mesh (10 296 elements, 10 566 nodes)

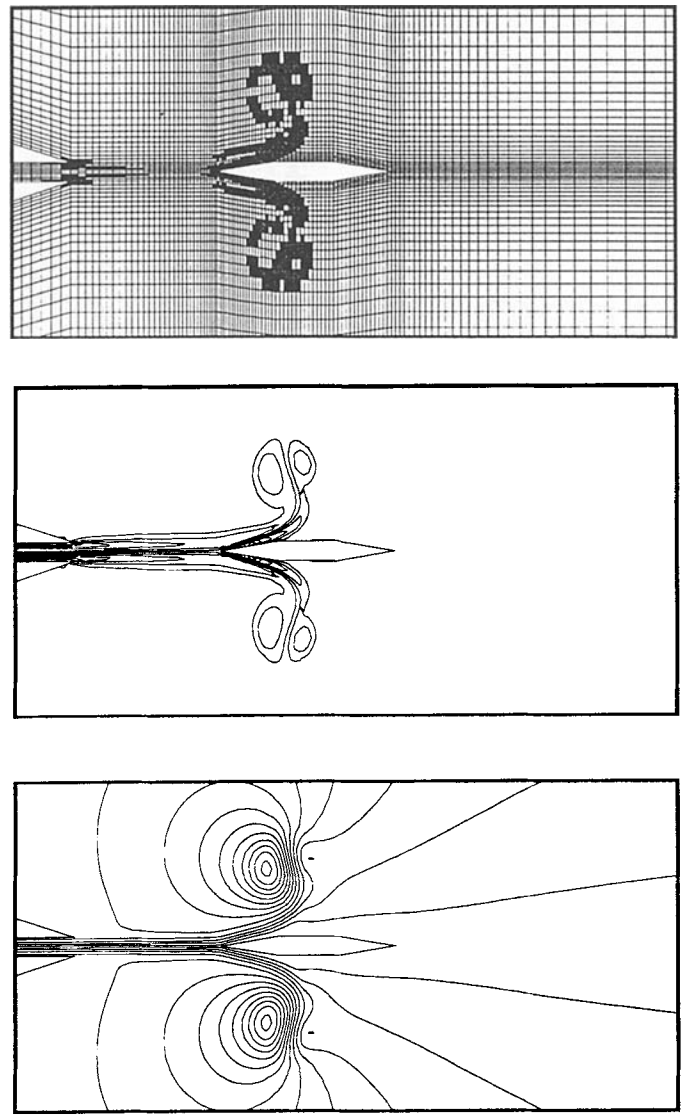

Figure 3. Plane jet impinging on a wedge at Reynolds number 250: solution obtained by the AIE/GEBE method at $t = 28{\cdot}75$. From top to bottom: distribution of the implicit elements, the vorticity and the streamlines

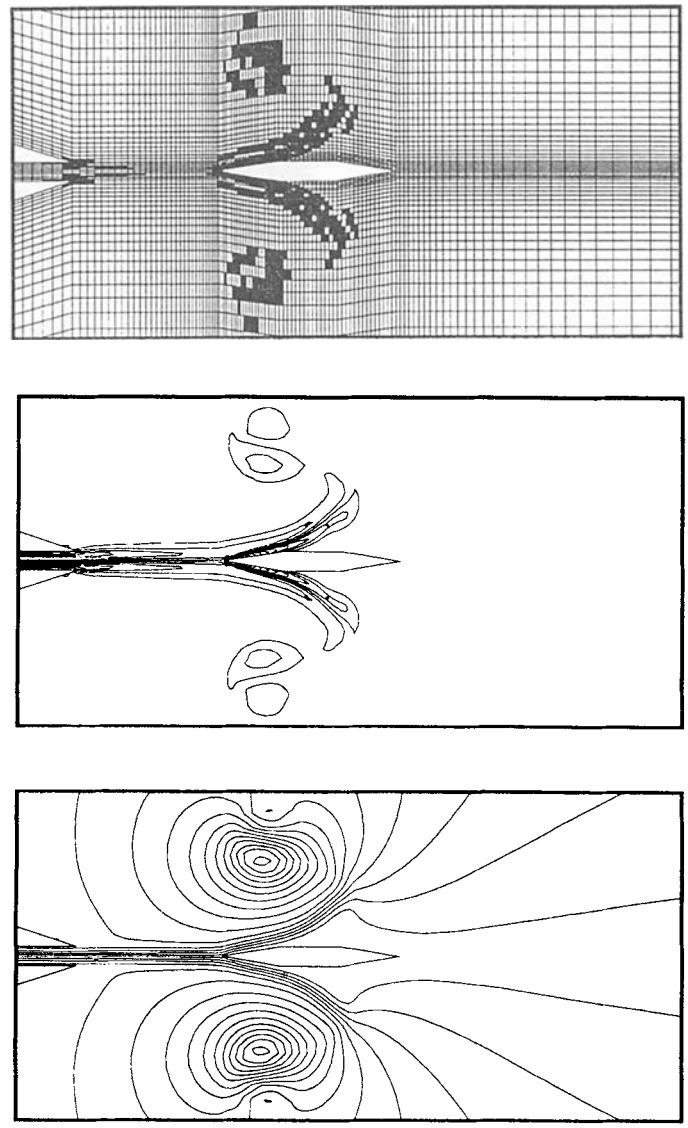

Figure 4. Plane jet impinging on a wedge at Reynolds number 250: solution obtained by the AIE/GEBE method at t = 37·50. From top to bottom: distribution of the implicit elements, the vorticity and the streamlines

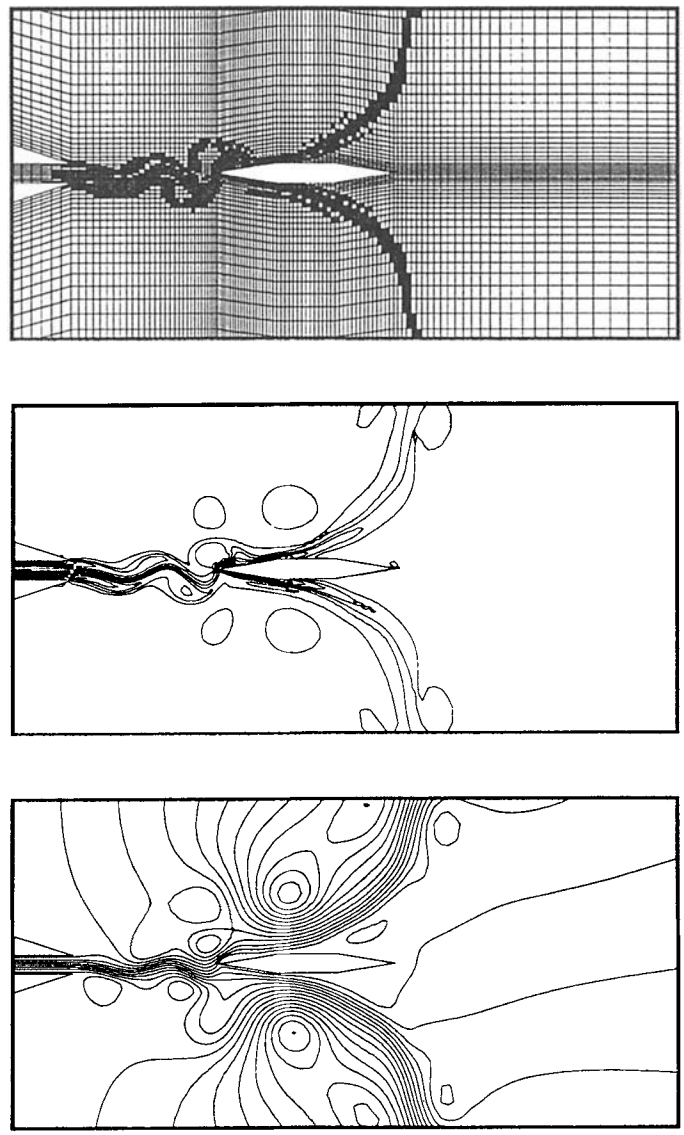

Figure 5. Plane jet impinging on a wedge at Reynolds number 250: solution obtained by the AIE/GEBE method at t = 63·75. From top to bottom: distribution of the implicit elements, the vorticity and the streamlines

convergence limit for the iterative solvers is 10^{-7} for block 1 and 10^{-6} for block 2. Figures 3–5 show, at various time steps, the distribution of the implicit elements, the vorticity and the streamlines.

5. CONCLUSIONS

We have presented the iterative versions of the adaptive implicit–explicit method. In this approach the GEBE and GMRES iteration methods are employed to solve the equation systems resulting from the implicitly treated elements, and therefore no direct solution effort is involved. The benchmarking results demonstrate that the methods presented can substantially reduce the CPU time and memory requirements in large-scale flow problems.

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